# Symmetric functions and interacting particle processes 

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## Séminaire Flajolet

February 2, 2023
joint with Arvind Ayyer and James Martin, arXiv:2011.06117, arXiv:2209.09859

## particle models and symmetric functions

(1) the asymmetric simple exclusion process (ASEP) $\rightarrow$ combinatorial formula for Macdonald polynomials and some nice specializations
(2) modified Macdonald polynomials $\rightarrow$ the multispecies totally asymmetric zero range process (mTAZRP) and observables

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- integrable systems: a class of dynamical systems with a certain restricted structure, in particular making them solvable
- we are interested in studying integrable systems whose exact solutions (e.g. stationary distributions) can be expressed in terms of combinatorial formulas or special functions (e.g. Macdonald polynomials)
- the field was initiated by Spitzer in his 1970 paper where he defined the ASEP (Asymmetric Simple Exclusion Process) and the ZRP (Zero Range Process)


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single species ASEP
multispecies ASEP
- dynamics: any two adjacent particles may swap with some predetermined rate (in our case, fixed by a parameter $0 \leq t \leq 1$ ):

$$
X A B Y \xrightarrow{1} X B A Y \quad \text { and } \quad X B A Y \xrightarrow{t} X A B Y \quad \text { for } \quad A>B
$$

## our setting: ASEP on a circle



$$
n=8, \quad \lambda=(3,2,2,2,1,0,0,0)
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- Fix a circular lattice on $n$ sites, and choose $n$ nonnegative integer weights. Collection of weights: a vector $\lambda=\left(\lambda_{1} \geq \cdots \geq \lambda_{n} \geq 0\right)$.


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- $\operatorname{ASEP}(\lambda)$ is a Markov chain whose states are the compositions $\alpha \in \operatorname{Sym}(\lambda)$ that are rearrangements of $\lambda$ (on a circle: $\alpha_{n+1}=\alpha_{1}$ )


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- For example, $\operatorname{ASEP}((2,2,1,0))$ has 12 states:

$$
(2,2,1,0),(2,1,2,0),(2,1,0,2),(2,2,0,1),(2,0,2,1),(2,0,1,2),(0,2,2,1), \cdots
$$

The transitions from state $(2,1,2,0)$ are:

- $(1,2,2,0)$ with probability $t / 4$
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- Goal: compute the stationary probabilities


## Example for $\lambda=(2,1)$ and $n=3$

- the elements in the state space are:
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- the (row stochastic) transition matrix is:

$$
\left(\begin{array}{cccccc}
1-\frac{2+t}{3} & \frac{1}{3} & \frac{1}{3} & 0 & 0 & \frac{t}{3} \\
\frac{t}{3} & 1-\frac{1+2 t}{3} & 0 & \frac{t}{3} & \frac{1}{3} & 0 \\
\frac{t}{3} & 0 & 1-\frac{1+2 t}{3} & \frac{1}{3} & \frac{t}{3} & 0 \\
0 & \frac{1}{3} & \frac{t}{3} & 1-\frac{2+t}{3} & 0 & \frac{1}{3} \\
0 & \frac{t}{3} & \frac{1}{3} & 0 & 1-\frac{2+t}{3} & \frac{1}{3} \\
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\end{array}\right)
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- the (unnormalized) stationary distribution is:

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\begin{gathered}
\widetilde{\operatorname{Pr}}((2,1,0))=\widetilde{\operatorname{Pr}}((1,0,2))=\widetilde{\operatorname{Pr}}((0,2,1))=2+t \\
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$$

## From ASEP to Macdonald polynomials

Define the partition function of $\operatorname{ASEP}(\lambda, n)$ :

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\mathcal{Z}_{\lambda, n}(t)=\sum_{\alpha \in S_{n} \cdot \lambda} \widetilde{\operatorname{Pr}}(\alpha)(t)
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## Theorem (Cantini-de Gier-Wheeler '15)

The partition function of $\operatorname{ASEP}(\lambda, n)$ is a specialization of the Macdonald polynomial:

$$
P_{\lambda}(1, \ldots, 1 ; 1, t)=\mathcal{Z}_{\lambda, n}(t)
$$

## Symmetric functions

- Let $X=x_{1}, x_{2}, \cdots$ be a family of indeterminates, and let $\Lambda=\Lambda_{\mathbb{Q}}$ be the algebra of symmetric functions in $X$ over $\mathbb{Q}$
- $f\left(x_{1}, \ldots, x_{n}\right) \in \Lambda$ is symmetric if $\forall \pi \in S_{n}, f\left(x_{1}, \ldots, x_{n}\right)=f\left(x_{\pi(1)}, \ldots, x_{\pi(n)}\right)$


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- $\Lambda$ has several nice bases: e.g. $\left\{m_{\lambda}\right\},\left\{e_{\lambda}\right\},\left\{h_{\lambda}\right\},\left\{p_{\lambda}\right\}$, indexed by partitions $\lambda$. E.g. $m_{(2,1)}=\sum_{i, j} x_{i}^{2} x_{j}^{1}=x_{1}^{2} x_{2}+x_{1}^{2} x_{3}+x_{2}^{2} x_{3}+x_{1} x_{2}^{2}+x_{1} x_{3}^{2}+x_{2} x_{3}^{2}+\cdots$

Let $\langle$,$\rangle be the standard inner product on \Lambda .\left\{s_{\lambda}\right\}$ is the unique basis of $\Lambda$ :
i. orthogonal with respect to $\langle$,
ii. upper triangular with respect to $\left\{m_{\lambda}\right\}$ :

$$
s_{\lambda}=m_{\lambda}+\sum_{\mu<\lambda} c_{\mu \lambda} m_{\mu}
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where $<$ is with respect to dominance order on partitions.

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- $s_{\lambda}=\sum_{\sigma} x^{\sigma}$ where $\sigma$ 's are semi-standard fillings of the Young diagram of shape $\lambda$
E.g. $s_{(2,1)}=x_{1}^{2} x_{2}+x_{1}^{2} x_{3}+x_{1} x_{2}^{2}+x_{1} x_{2} x_{3}+x_{1} x_{2} x_{3}+x_{1} x_{3}^{2}+x_{2}^{2} x_{3}+x_{2} x_{3}^{2}=m_{(2,1)}+m_{(1,1,1)}$

| 2 |  | 3 |  | 2 |  | 3 |  | 2 |  | 3 |  | 3 |  | 3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | 1 | 1 | 1 | 2 | 1 | 2 | 1 | 3 | 1 | 3 | 2 | 2 | 2 | 3 |

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Let $\langle,\rangle_{q, t}$ be the inner product on $\Lambda(q, t)$. Then $\left\{P_{\lambda}\right\}$ is the unique basis of $\Lambda(q, t)$ that is uniquely determined by:
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P_{\lambda}(X ; q, t)=m_{\lambda}(X)+\sum_{\mu<\lambda} c_{\mu \lambda}(q, t) m_{\mu}(X)
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- Example:

$$
P_{(2,1)}(X ; \boldsymbol{q}, t)=m_{(2,1)}+\frac{(1-t)(2+q+t+2 q t)}{1-q t^{2}} m_{(1,1,1)}
$$

## Combinatorial formulas

- Haglund-Haiman-Loehr '04 gave a formula for $P_{\lambda}$ as a sum over tableaux with statistics maj and (co)inv:

$$
P_{\lambda}(X ; q, t)=\sum_{\substack{\sigma \in \operatorname{dg}(\lambda) \\ \sigma \text { non-attacking }}} x^{\sigma} q^{\operatorname{maj}(\sigma)} t^{\operatorname{coinv}(\sigma)} \prod_{u} \frac{1-t}{1-q^{\operatorname{leg}(u)+1} t^{\operatorname{arm}(u)+1}}
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- Corteel-M-Williams '18: a new formula for $P_{\lambda}$ in terms of multiline queues, which also give formulas for the stationary distribution of the ASEP; this was inspired by the result of Cantini-de Gier-Wheeler '15


## multiline queues

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Corteel-M-Williams '18 (general)

- The state of a multiline queue is read off the bottom row
- The weight $w t(M)$ of a multiline queue depends on the parameters $t, q, x_{1}, \ldots, x_{n}$ :

$$
\text { weight }=x_{1}^{2} x_{2}^{2} x_{3} x_{4}^{2} x_{5} x_{6}^{2} q t^{2} \frac{(1-t)^{3}}{\left(1-q t^{3}\right)^{2}\left(1-q t^{2}\right)}
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| 2 | 4 |  |  |  |
| :--- | :--- | :--- | :--- | :--- |
| 6 | 1 | 3 |  |  |
| 6 | 1 | 5 | 2 | 3 |

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$$
\begin{aligned}
n & =6 \\
\lambda & =(3,3,2,1,1) \\
\lambda^{\prime} & =(2,3,5)
\end{aligned}
$$

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| :--- | :--- | :--- | :--- | :--- |
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- Can be represented by a tableau, where each string is mapped to a column
- Can be represented by a queueing system and described as a coupled system of 1-ASEPs. The pairing is a projection map onto the $n$-ASEP.


## From ASEP to Macdonald polynomials

Theorem (Martin '18, Corteel-M-Williams '18)
The (unnormalized) stationary probability of state $\alpha$ of the mASEP is

$$
\widetilde{\operatorname{Pr}}(\alpha)(t)=\sum_{M \in M L Q(\alpha)} w t(M)(1, \ldots, 1 ; 1, t)
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## Theorem (Cantini-de Gier-Wheeler '15)

The partition function of $\operatorname{ASEP}(\lambda, n)$ is a specialization of the Macdonald polynomial:

$$
P_{\lambda}(1, \ldots, 1 ; 1, t)=\mathcal{Z}_{\lambda, n}(t)=\sum_{\alpha \in S_{n} \cdot \lambda} \widetilde{\operatorname{Pr}}(\alpha)(t)
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$$
P_{\lambda}\left(x_{1}, \ldots, x_{n} ; q, t\right)=\sum_{M \in \operatorname{MLQ}(\lambda, n)} w t(M)\left(x_{1}, \ldots, x_{n} ; q, t\right)
$$

(also Lenart '09 for $\lambda$ with distinct parts.)

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- at $t=0$, the ASEP becomes the TASEP:

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- In the multiline queues, pairings become forced.


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- In the multiline queues, pairings become forced.

|  | (3) |  | (3) |  |  | $n=6$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| (2) |  | (3) |  |  | (3) | $\lambda=(3,3,2,1,1)$ |
| (2) | (1) |  | (3) | (1) | (3) | $\alpha=(2,1,0,3,1,3)$ |
| 1 | 2 | 3 | 4 | 5 | 6 |  |

- call $\operatorname{MLQ}(\lambda, n)$ the set of ball arrangements with $\lambda_{j}^{\prime}$ balls in each row $j$, on a lattice of size $n \times \lambda_{1}$. (The labels can be recovered uniquely)

$$
M=(\{1,2,4,5,6\},\{1,3,6\},\{2,4\})
$$

## $q=t=0$ : Schur polynomials via multiline queues

- at $q=0$, the multiline queues are non-wrapping, denote this set by $\operatorname{MLQ}_{0}(\lambda, n)$ :


| $\square$ | $\bullet \bullet$ |
| :--- | :--- | :--- |
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$$
\sum_{M \in \operatorname{MLQ}_{0}(\lambda, n)} x^{M}=s_{\lambda}
$$




$$
\begin{array}{|l|l|}
\hline 4 & 4 \\
\hline 2 & 3 \\
\hline
\end{array} \quad \begin{array}{|l|l|}
\hline 4 & 4 \\
\hline 3 & 3 \\
\hline
\end{array}
$$

- the map $\operatorname{MLQ}_{0}(\lambda) \rightarrow \operatorname{SSYT}(\lambda)$ is given by column $\operatorname{RSK}$ applied to the row reading word of the multiline queue. (bottom to top, left to right)


## Lascoux-Schützenberger charge formula via MLQs

- For a permutation $\sigma$, define charge $(\sigma):=\operatorname{maj}\left(\operatorname{rev}\left(\sigma^{-1}\right)\right)$. For a SSYT $\tau$, charge $(\tau)=\operatorname{charge}(\operatorname{rw}(\tau))$.


## Theorem (Lascoux-Schützenberger, '78)

$$
P_{\lambda}(X ; q, 0)=\sum_{\mu \leq \lambda} K_{\mu^{\prime} \lambda^{\prime}}(q) s_{\mu}, \quad K_{\lambda \mu}(q)=\sum_{Q \in \operatorname{SSYT}(\lambda, \mu)} q^{\text {charge }(Q)}
$$

- Define collapsing procedure $\rho$ (with Jerónimo Valencia '23+):

$$
\begin{array}{rr}
\operatorname{MLQ}(\lambda) & \bigcup_{\mu} \operatorname{MLQ}_{0}(\mu) \times \operatorname{SSYT}\left(\mu^{\prime}, \lambda^{\prime}\right) \\
M \longrightarrow & \left(M_{0}, Q\right)
\end{array}
$$

- charge $(Q)=\operatorname{maj}(M)$ (the $q$-statistic, keeps track of wrapping pairings)
- can be described using lowering operators on the column reading word of $M$
- lifting procedure $\rho^{-1}$ can be described using raising operators
- generalizes to a quasisymmetric refinement of $K_{\lambda \mu}(q)$.


## Collapsing procedure



$$
\begin{gathered}
M \in \operatorname{MLQ}((6,4,2)) \\
\operatorname{maj}(M)=2+1+3=6
\end{gathered}
$$

## Collapsing procedure



## Collapsing procedure



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## Collapsing procedure




$$
\begin{aligned}
& \lambda=(6,4,2) \\
& \operatorname{maj}(M)=6
\end{aligned}
$$



$$
\begin{gathered}
\lambda=(4,3,2,2,1) \\
\operatorname{maj}\left(M_{0}\right)=0
\end{gathered}
$$

| 4 |  |  |  |  |
| :--- | :--- | :--- | :--- | :---: |
| 3 | 6 |  |  |  |
| 2 | 2 | 3 | 4 |  |
|  |  |  |  |  |
| 1 | 1 | 1 | 2 |  |
| 5 |  |  |  |  |

charge $(Q)=6$

## Collapsing procedure

## Lemma

$$
\begin{aligned}
\operatorname{RSK}^{c o l}:\binom{a_{1}, \ldots, a_{n}}{b_{1}, \ldots, b_{n}} & \rightarrow \operatorname{SSYT}\left(\mu^{\prime}\right) \times \operatorname{SSYT}(\mu) \\
\operatorname{RSK}^{c o l}(\operatorname{rw}(M)) & =\operatorname{RSK}^{c o l}(\operatorname{rw}(\rho(M))
\end{aligned}
$$


$r w(M)=234|135| 23|15| 6 \mid 2$
$\operatorname{rw}(\rho(M))=23456|1235| 13 \mid 2$

## Littlewood-Richardson rule via MLQs

## Theorem

$$
s_{\lambda} s_{\mu}=\sum_{\nu} c_{\lambda \mu}^{\nu} s_{\nu}=\sum_{\nu} c_{\lambda^{\prime} \mu^{\prime}}^{\nu^{\prime}} s_{\nu}
$$

- Using the collapsing procedure:

$$
\operatorname{MLQ}_{0}(\lambda) \times \operatorname{MLQ}_{0}(\mu) \rightarrow \operatorname{MLQ}_{0}(\nu) \times \operatorname{SSYT}^{*}\left(\nu^{\prime} / \lambda^{\prime}, \mu\right)
$$



- $\operatorname{SSYT}^{*}\left(\nu^{\prime} / \lambda^{\prime}, \mu^{\prime}\right)$ is the set of Yamanouchi fillings of $\nu^{\prime} / \lambda^{\prime}$ with content $\mu^{\prime}$


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## modified Macdonald polynomials $\widetilde{H}_{\lambda}(X ; q, t)$

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- Corteel-Haglund-M-Mason-Williams '20 conjecture: a new formula for $\widetilde{H}_{\lambda}$ with maj and a new statistic quinv:

$$
\widetilde{H}_{\lambda}(X ; q, t)=\sum_{\sigma \in \operatorname{dg}(\lambda)} q^{\operatorname{maj}(\sigma)} t^{q u i n v(\sigma)} x^{\sigma}
$$

## From multiline queues to a new formula for $\widetilde{H}_{\lambda}$

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- this leads to a new "queue inversion" statistic for $t$ that we call quinv:



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- this leads to a new "queue inversion" statistic for $t$ that we call quinv:

(Corteel-Haglund-M-Mason-Williams '20, Ayyer-M-Martin '21)
- the resulting objects are of the same flavor as multiline queues, except that multiple balls are allowed at each location. (This translates to removing the "non-attacking" condition from the corresponding tableaux)


## Example: $\tilde{H}_{(2,1)}(X ; q, t)$

$\widetilde{H}_{(2,1)}\left(x_{1}, x_{2} ; q, t\right)=m_{(3)}+(1+t+q) m_{(2,1)}+(1+2 t+2 q+q t) m_{(1,1,1)}$

- (AMM21) $\quad \widetilde{H}_{\lambda}(X ; q, t)=\sum_{\sigma: \mathrm{dg}(\lambda) \rightarrow \mathbb{Z}_{+}} q^{\text {maj }(\sigma)} t^{\text {quinv }(\sigma)} x^{\sigma}$

| 1 |  | 2 |  | 1 |  | 1 |  | 1 |  | 2 |  | 3 |  | 2 |  | 1 |  | 3 |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | 1 | 1 | 1 | 2 | 2 | 1 | 2 | 3 | 1 | 3 | 1 | 2 | 3 | 1 | 3 | 2 | 2 | 1 |

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| 1 | 2 | 1 | 1 | 1 | 2 | 3 | 2 | 1 | 3 |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 11 | 11 | 12 | 21 | 23 | 13 | 12 | 31 | 32 | 21 |
| $m_{3}$ | $q m_{21}$ | $m_{21}$ | $t m_{21}$ | $m_{111}$ | $q m_{111}$ | q $m_{111}$ | $t m_{111}$ | $t m_{111}$ | qt $m_{111}$ |

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## motivation for queue inversions: multiline diagrams

The tableaux are actually representing a queueing system which is an arrangement of lattice paths/strings


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"plethystic version" of certain non-attacking fillings queue inversion

"plethystic version" of multiline queues
skipped particle

## Big picture



What is the analogous interacting particle system whose partition function is a specialization of $\widetilde{H}_{\lambda}$ ?

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## totally asymmetric zero range processes (TAZRP)

- continuous-time stochastic processes (Spitzer '70), can be defined on arbitrary graphs. In our case, we have a circular lattice with $n$ sites.



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$$
\begin{gathered}
\text { Here, } n=5, k=7 \\
\tau=(11|\cdot| 111|1| 1)
\end{gathered}
$$

- simplest case: there are $k$ indistinguishable particles, moving clockwise. A configuration $\tau=\left(\tau_{1}, \ldots, \tau_{n}\right)$ is any allocation of the $k$ particles on the $n$ sites.


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- multispecies variant: we now allow different particle types, labeled by integers (particles of the same type are still indistinguishable)
- Kuniba-Maruyama-Okado (2015+) (and others) have studied many multispecies variants of the TAZRP. All of these are integrable! The version we will describe was first studied by Takayama '15


## the mTAZRP: states

- Fix a (circular 1D) lattice on $n$ sites and a partition $\lambda=\left(\lambda_{1} \geq \cdots \geq \lambda_{k}>0\right)$ for the particle types



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- Fix a (circular 1D) lattice on $n$ sites and a partition $\lambda=\left(\lambda_{1} \geq \cdots \geq \lambda_{k}>0\right)$ for the particle types
- $\operatorname{TAZRP}(\lambda, n)$ is a Markov chain whose states are multiset compositions $\tau$ of type $\lambda$, with $n$ (possibly empty) parts



## the mTAZRP: transition rates

- Each particle is equipped with an exponential clock. Transitions are jumps from site $j$ to site $j+1$


## the mTAZRP: transition rates

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$$

For example: If site $j$ contains the particles $\{4,3,3,1,1,1\}$, then:

$$
\begin{array}{lll}
k=1: & d=3, & c=3, \\
k=3: & d=1, & c=2, \\
k=4: & d=0, & c=1,
\end{array}
$$

## Lumping of tableaux to mTAZRP

Very similar projection map as for the ASEP.

- Given a filling $\sigma$, read the state $\tau \in \operatorname{TAZRP}(\lambda, n)$ from the bottom row of $\sigma$ as follows:
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$\tau_{j}$ is the multiset $\left\{\lambda_{i}: \sigma(1, i)=j\right\}$
- For example, for $\lambda=(2,1,1)$ and $n=3$, the following are all the tableaux that correspond to the state $\tau=(21|\cdot| 1)$ :

| 1 |  |  | 2 |  | 3 |  |  | 1 |  |  | 2 |  |  | 3 |  |  |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 1 | 1 | 3 | 1 | 13 | 1 | 1 | 3 | 1 | 3 | 1 | 1 | 3 | 1 | 1 | 3 |  |

## TAZRP probabilities and tableaux

## Theorem (Ayyer-M-Martin '21)

Fix $\lambda, n$. The (unnormalized) stationary probability of $\tau \in \operatorname{TAZRP}(\lambda, n)$ is

$$
\tilde{\operatorname{Pr}}(\tau)=\sum_{\substack{\sigma: \operatorname{dg}(\lambda) \rightarrow[n] \\ \sigma \text { has type } \tau}} x^{\sigma} t^{\text {quinv }(\sigma)}
$$

## Corollary

The so-called partition function of $\operatorname{TAZRP}(\lambda, n)$ is

$$
\mathcal{Z}_{\lambda, n}\left(x_{1}, \ldots, x_{n} ; t\right)=\widetilde{H}_{\lambda}\left(x_{1}, \ldots, x_{n} ; 1, t\right)
$$

Proof: construction of a Markov chain on tableaux that lumps to the TAZRP.

## an example for $\lambda=(2,1,1)$ and $n=2$

The stationary distribution is:

Example computation for $(21 \mid 1)$ :

| 1 |  |  |
| :--- | :--- | :--- |
| 1 | 1 | 2 |\(: \begin{aligned} \& <br>

\& t^{2},\end{aligned}\)

$$
\begin{array}{|l|l|l}
\hline 2 & & \\
\hline 1 & 1 & 2 \\
\hline
\end{array}
$$

$$
\begin{array}{|l|l|l}
\hline 1 & & \\
\hline 1 & 2 & 1 \\
\hline
\end{array}
$$

$$
\begin{array}{|l|l|l}
\hline 2 & & \\
\hline 1 & 2 & 1 \\
\hline
\end{array}
$$

the total is:

$$
\widetilde{\operatorname{Pr}}(21 \mid 1)=x_{1}^{2} x_{2}\left(t x_{1}+x_{2}\right)(1+t) .
$$

## Current

- The current of particle $\ell$ across the edge $j$ is defined as the number of particles of type $\ell$ traversing the edge $j$ per unit of time in the large time limit.


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$$
\begin{gathered}
\text { Here, } n=5, m=7 \\
\tau=(2,0,3,1,1)
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- The stationary probability of the configuration $\tau$ is:

$$
\pi(\tau)=\frac{1}{\widetilde{H}_{\left(1^{m}\right)}\left(x_{1}, \ldots, x_{n} ; 1, t\right)}\left[\begin{array}{c}
m \\
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$$

## Proposition (Current for the single species TAZRP)

For the single-species TAZRP on $n$ sites with $m$ particles, the current is given by

$$
J=[m]_{t} \frac{\widetilde{H}_{\left(1^{m-1}\right)}\left(x_{1}, \ldots, x_{n} ; 1, t\right)}{\widetilde{H}_{\left(1^{m}\right)}\left(x_{1}, \ldots, x_{n} ; 1, t\right)}
$$

## Current

## Theorem (Ayyer-M-Martin '22+)

Let $\lambda=\left(1^{m_{1}}, \ldots, k^{m_{k}}\right)$, and let $1 \leq j \leq k$. The current of the particle of type $j$ of the TAZRP of type $\lambda$ on $n$ sites is given by

$$
J=\left[m_{j}+\cdots+m_{k}\right]_{t} \frac{\widetilde{H}_{\left(1^{m_{j}+\cdots+m_{k}-1}\right)}}{\widetilde{H}_{\left(1^{m_{j}+\cdots+m_{k}}\right)}}-\left[m_{j+1}+\cdots+m_{k}\right]_{t} \frac{\widetilde{H}_{\left(1^{m_{j+1}+\cdots+m_{k}-1}\right)}}{\widetilde{H}_{\left(1^{m_{j+1}+\cdots+m_{k}}\right)}}
$$

## Densities

- Take $\operatorname{TAZRP}(\lambda, n)$ with content $\lambda=\left(1^{m_{1}}, 2^{m_{2}}, \ldots, k^{m_{k}}\right)$.
- Define $z_{j}^{(\ell)}$ to be the random variable counting the number of particles of type $\ell$ at site $j$ in a configuration of $\operatorname{TAZRP}(\lambda, n)$.
- Denote the expectation in the stationary distribution by $\left\langle z_{j}^{(\ell)}\right\rangle$.


## Theorem (Ayyer-M-Martin '22+)

For $1 \leq \ell \leq k$, the density of the $\ell$ 'th species at site 1 is given by

$$
\left\langle z_{1}^{(\ell)}\right\rangle=x_{1} \partial_{x_{1}} \log \left(\frac{\widetilde{H}_{\left(1^{m}{ }^{m}+\cdots+m_{k}\right)}\left(x_{1}, \ldots, x_{n} ; 1, t\right)}{\widetilde{H}_{\left(1^{m_{\ell+1}+\cdots+m_{k}}\right)}\left(x_{1}, \ldots, x_{n} ; 1, t\right)}\right) .
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## Corollary

$\left\langle z_{1}^{(\ell)}\right\rangle$ is symmetric in the variables $\left\{x_{2}, \ldots, x_{n}\right\}$.

## Local correlations

- Fix $\lambda, n$, and $0 \leq \ell \leq n$, and let $w$ be a configuration of the TAZRP on the first $\ell$ sites of type $\mu$, where $\mu \subseteq \lambda$.


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Configurations contributing to $\mathbb{P}_{\lambda, n}(\bar{w})$ are

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## Theorem (Ayyer-M-Martin '22)

$\mathbb{P}_{\lambda, n}(\bar{w})$ is tcbsymmetric in the variables $\left\{x_{\ell+1}, \ldots, x_{n}\right\}$.

## final remarks

- Explicit bijection from the inv to the quinv statistic?

(recently found by Loehr)
- Can we find a dynamical process that incorporates the $q$ as a parameter?

This seems difficult because

- We lose factorization of $\tilde{H}_{\lambda}$
- We lose translation invariance
- Suitable quasisymmetric version of modified Macdonald polynomials?

Nonsymmetric version?


Modified Macdonald polynomials and the multispecies zero range process: arXiv:2011.06117, arXiv:2209.09859

